

# Friction and adhesion in suspension rheology by constraint counting

James A. Richards, Ben M. Guy, Daniel Hodgson, Elena Blanco & Wilson C. K. Poon.

Contact: james.a.richards@ed.ac.uk

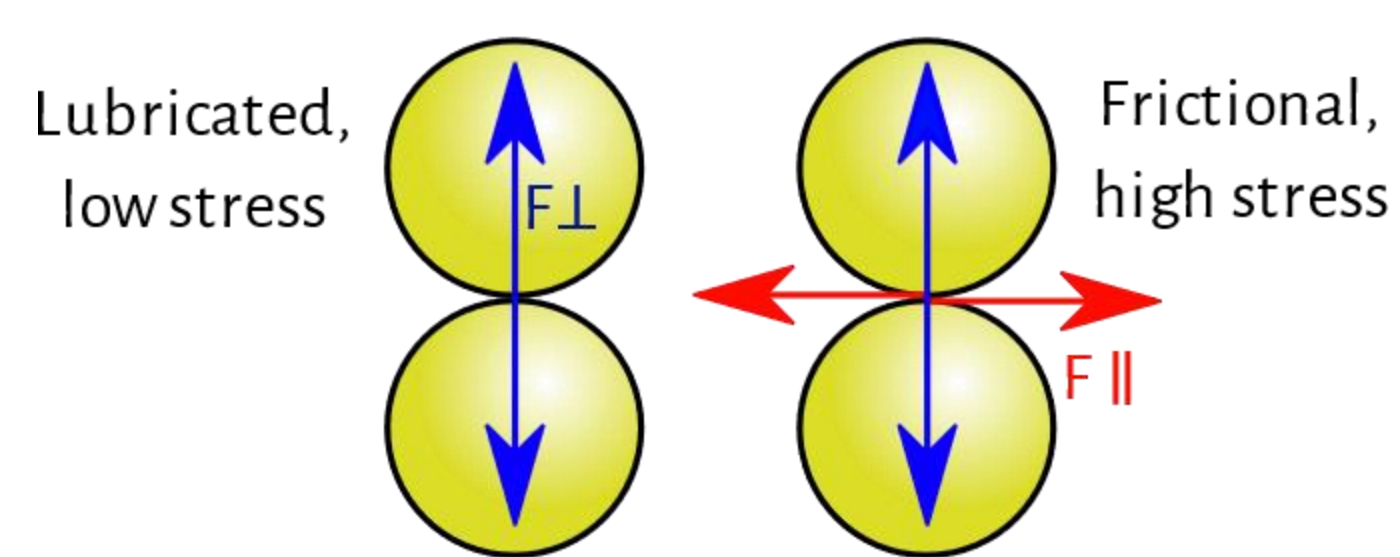


THE UNIVERSITY  
of EDINBURGH

**In theory:** suspensions of hard non-Brownian particles should have a constant viscosity. **In practice:** they are highly non-Newtonian; there must be additional relevant stress scales in these systems.

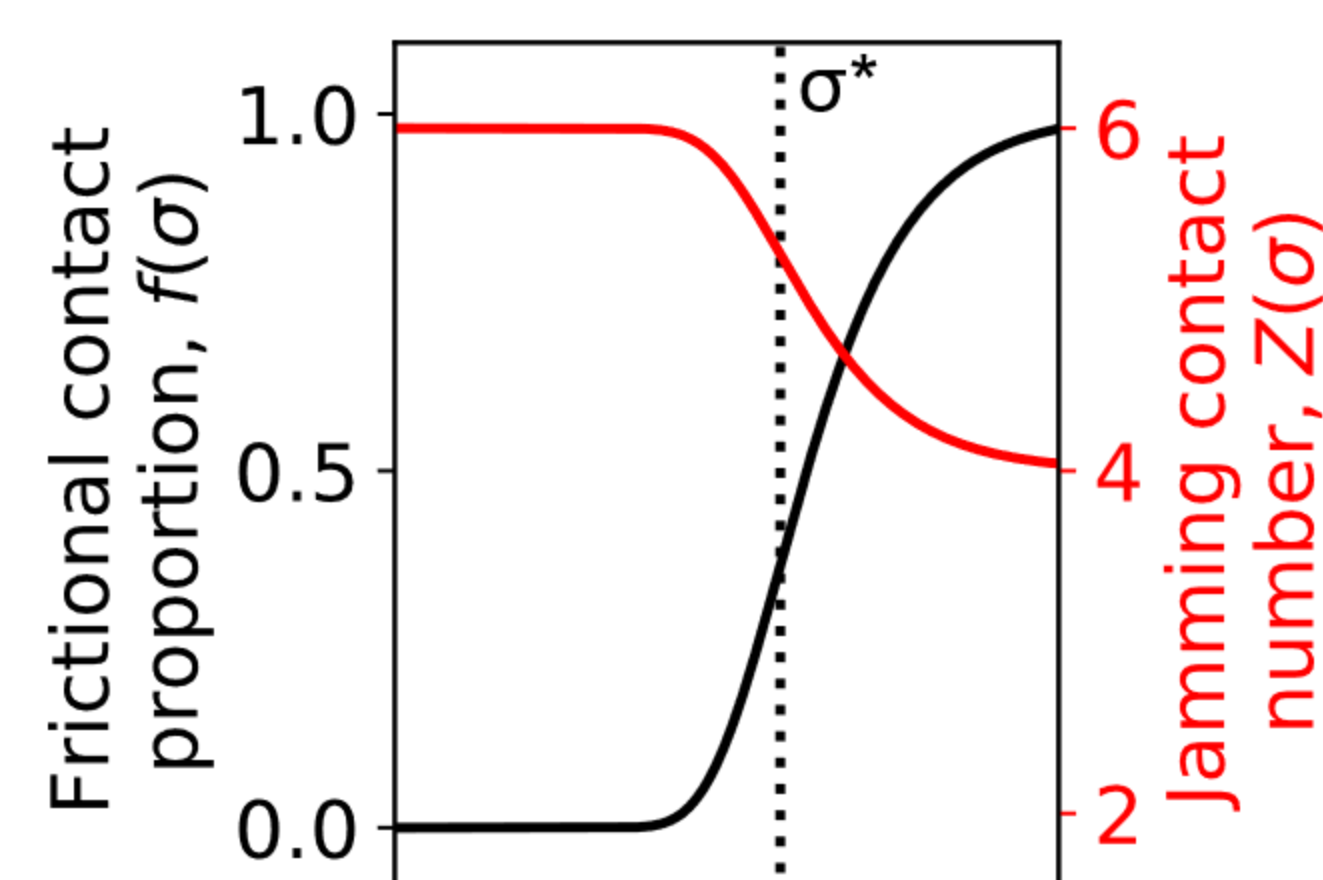
## Rheology by constraints: revisiting shear thickening and friction

Shear-thickening is driven by the onset of frictional contact at a **critical stress,  $\sigma^*$** , when the stabilisation is overcome.



### Recasting Wyart & Cates theory [2]

Stress sets the proportion of frictional contacts,  $f(\sigma)$ :  $f = \exp(-(\sigma/\sigma^*)^\beta)$

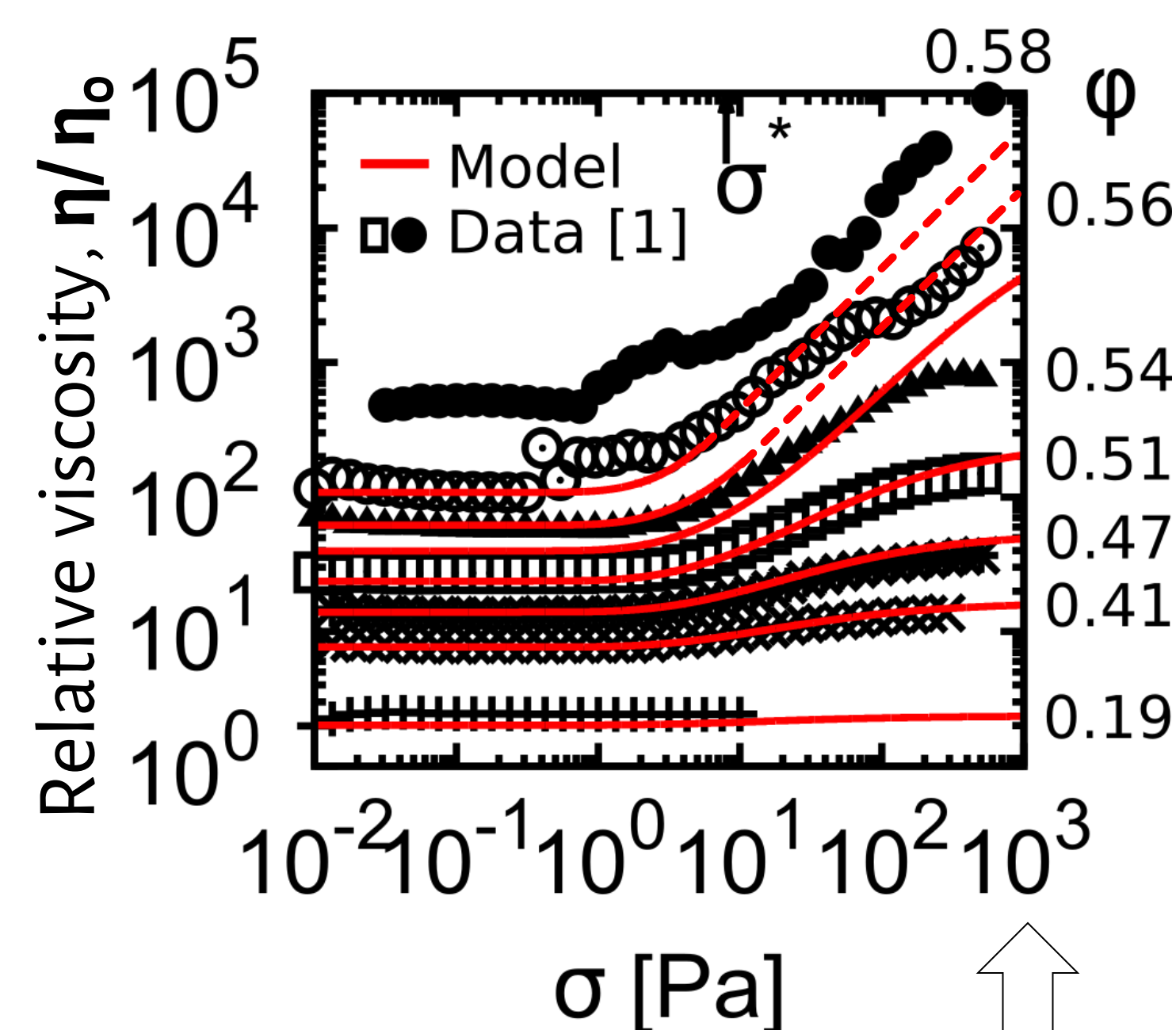


$f$  determines contacts needed to jam,  $Z(\sigma)$ :  $Z = 6 - 2f$

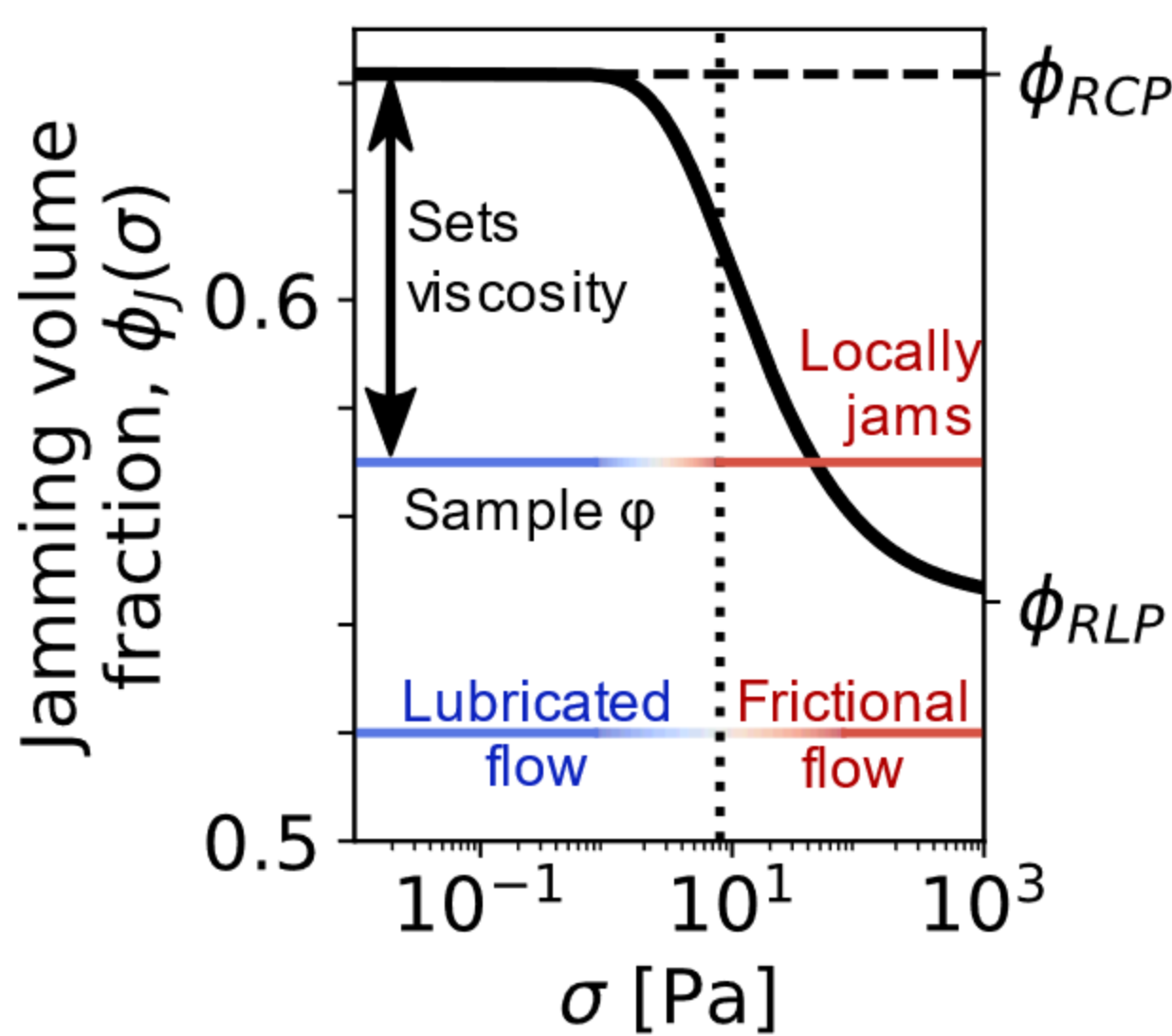
$Z$  controls the jamming volume fraction [3],  $\phi_J(\sigma)$ :

$$\phi_J = \frac{Z}{Z + C}$$

Use Kreiger-Dougherty to calculate viscosity,  $\eta(\sigma, \phi)$



**System:** 4  $\mu\text{m}$  sterically stabilised PMMA in CHB/decalin.  $C=3.35$  &  $\beta=0.8$



$$\eta = \frac{\eta_0}{(1 - \phi/\phi_J)^2}$$

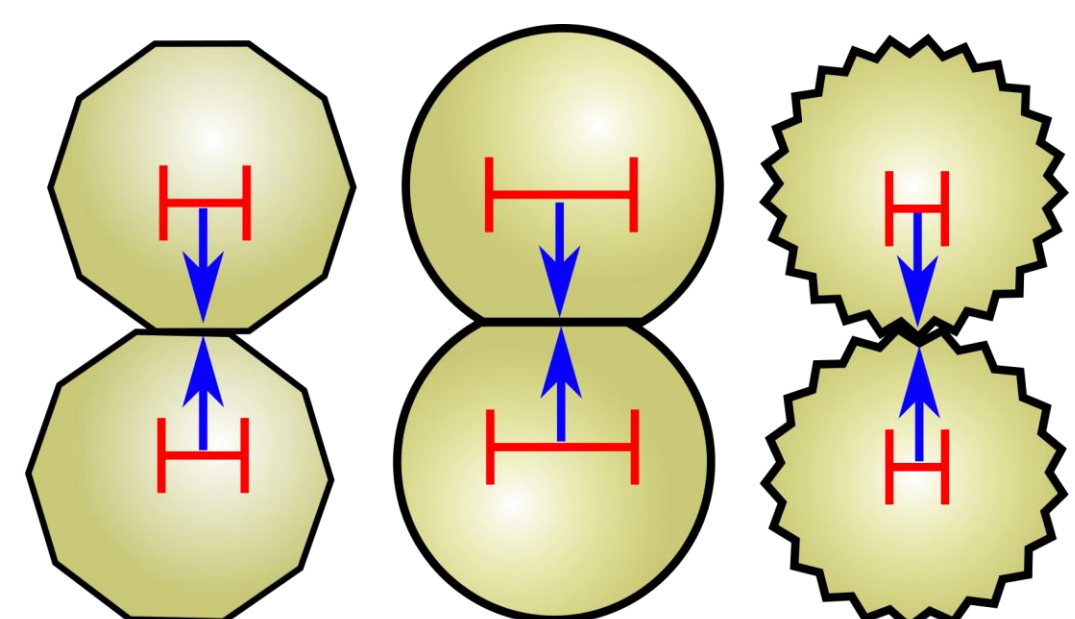
Model flow curves,  $\eta(\sigma, \phi)$

## Introducing another constraint: adhesion

Radial attraction does not add constraints; adhesion, **resisting rolling**, is the relevant effect. The strength will depend on both the interaction and on surface shape.

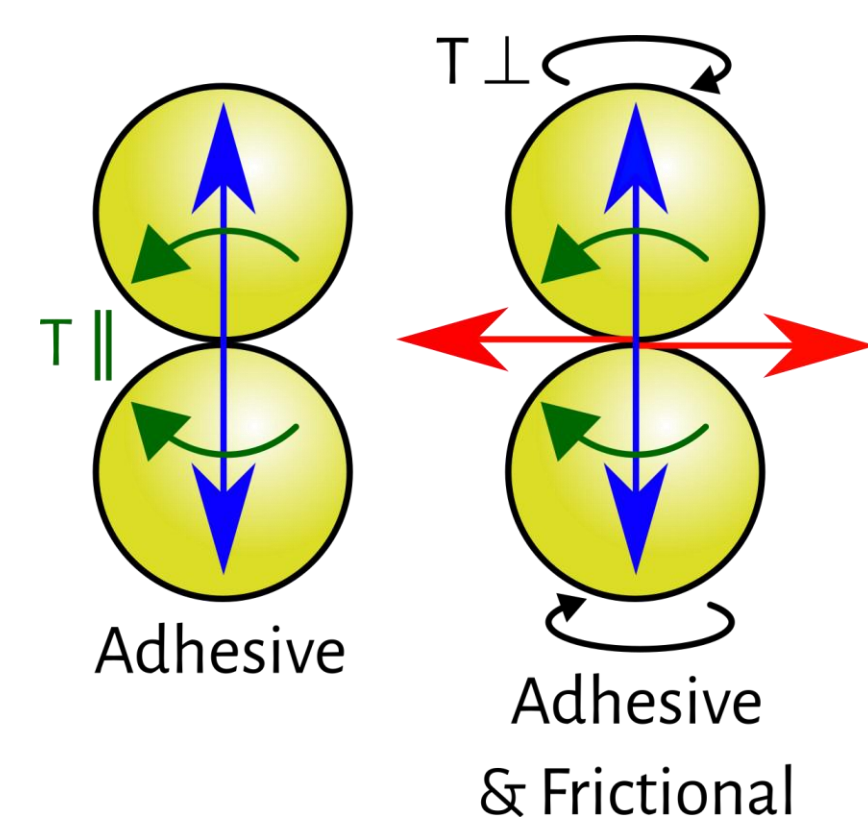
It requires a contact area, which may arise from:

- shape
- deformability, or
- roughness



The proportion of adhesive bonds  $a$  will **decrease** from 1 to 0 with stress above  $\sigma_A$ , the strength of adhesion:

$$a(\sigma) = 1 - \exp(-(\sigma_A/\sigma)^\kappa)$$



Consider complex interactions only in terms of simple constraints, now:

$$Z = 6 - 2a - 2f$$

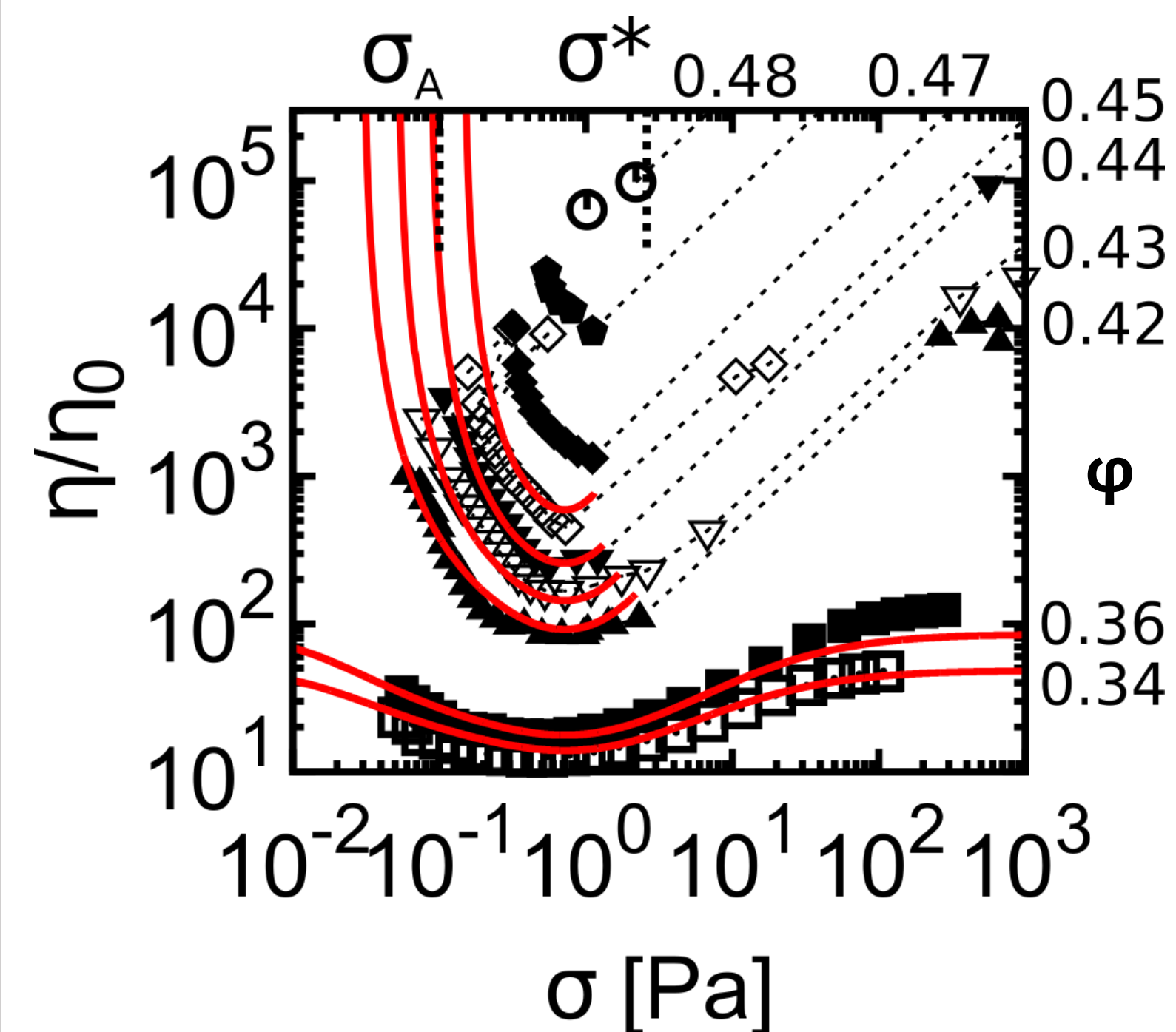
## References & Funding Acknowledgements

1. Guy, Hermes & Poon, *PRL* (2015). 2. Wyart & Cates, *PRL* (2014). 3. Song, Wang & Makse, *Nature* (2008).



**Conclusion:** many flow curves are captured by considering constraints (sliding, rolling, ...) at contacts. Rich behaviour arises from the balance of breaking adhesive bonds and making frictional contacts with changing stress.

## I. Weak adhesion: $\sigma_A \ll \sigma^*$

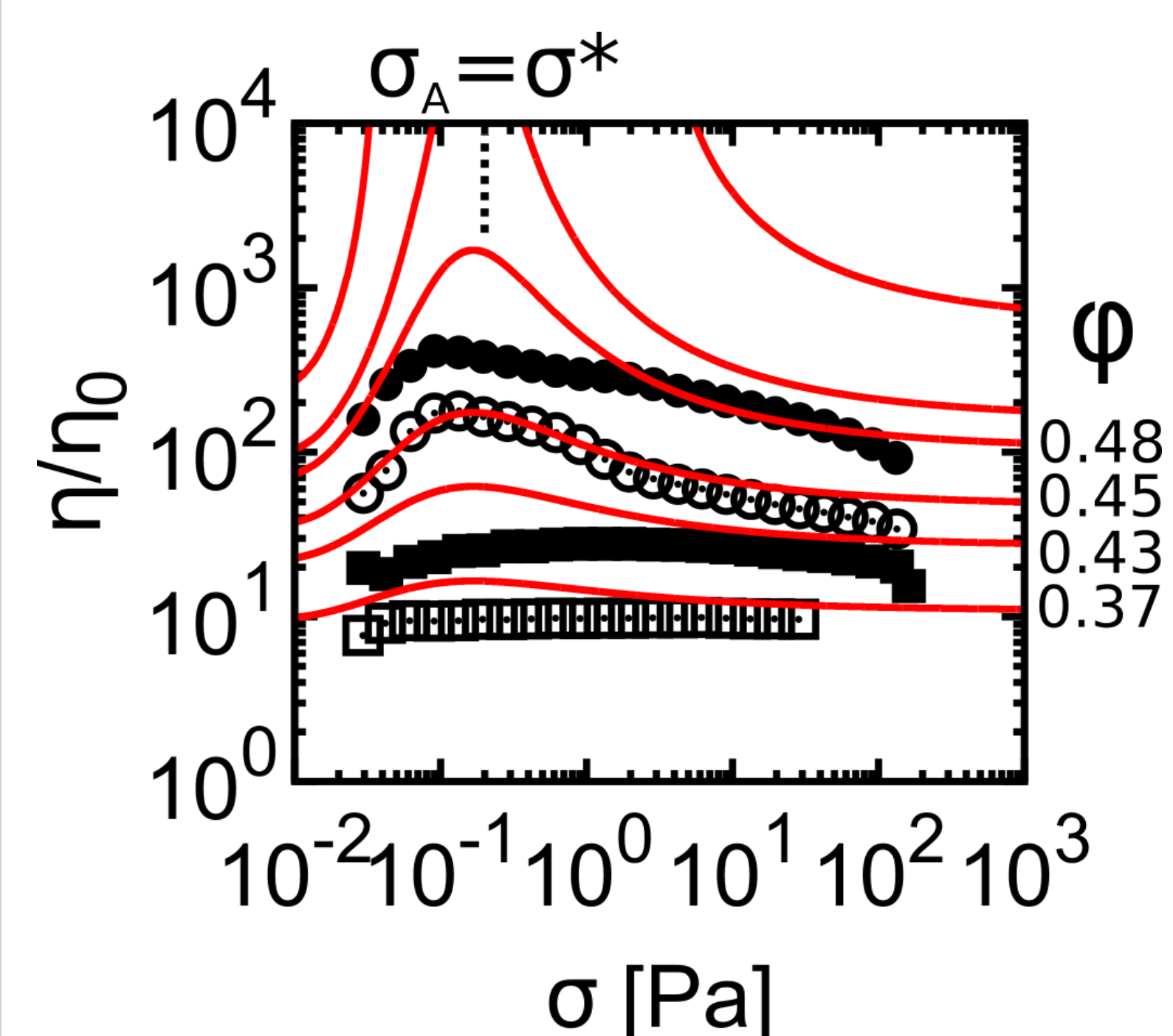


**System:** corn-starch in 50wt% water-glycerol.  $\beta=0.8$ ,  $\kappa=0.5$  &  $C=6$ .

The adhesive bonds are broken before frictional contacts are formed.

Both adhesion and friction introduce two constraints. A yield stress arises at the same  $\phi$  as Discontinuous Shear Thickening (slope 1 on figure).

## II. Intermediate adhesion: $\sigma_A \approx \sigma^*$

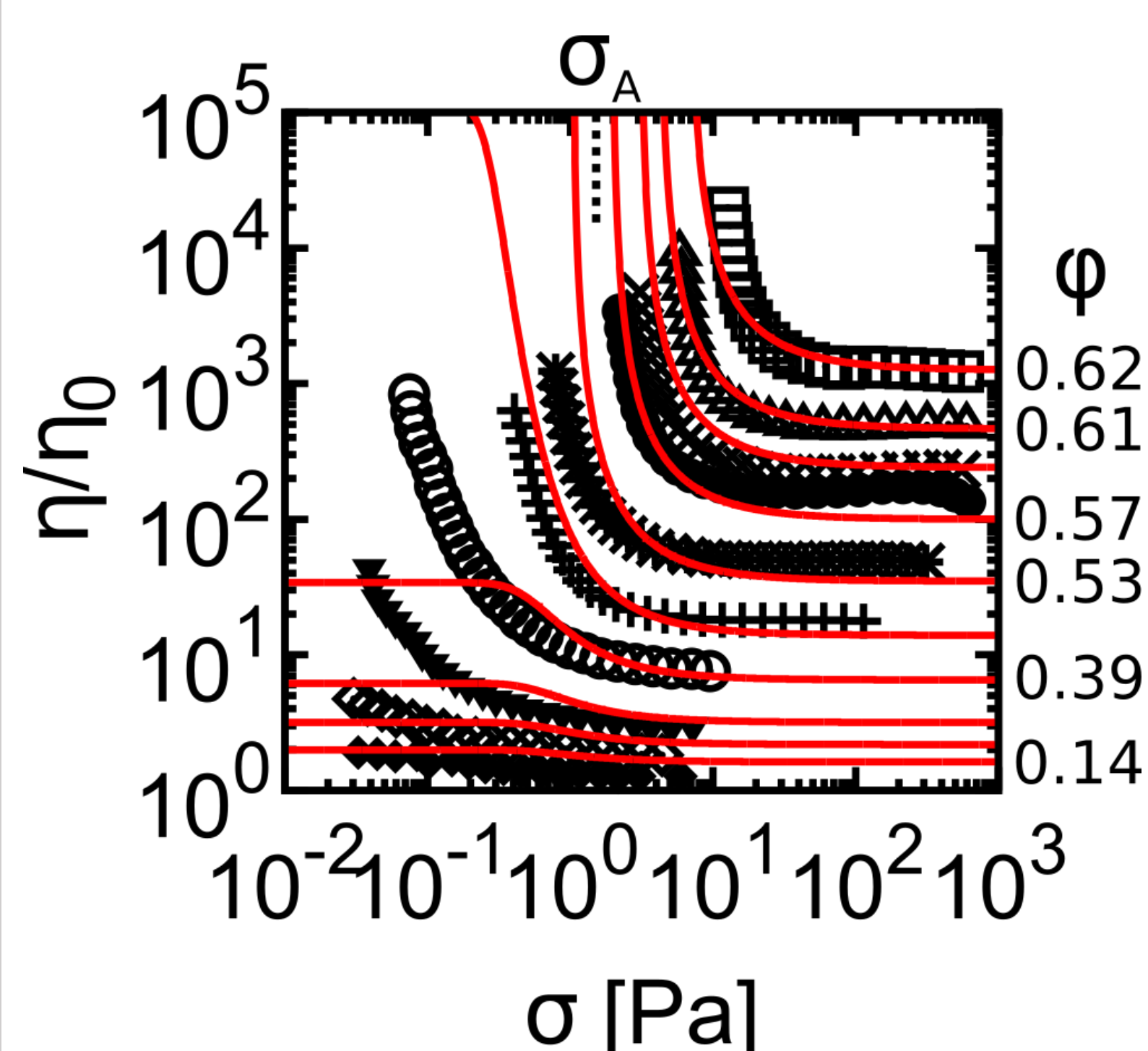


**System:** 45  $\mu\text{m}$  sterically stabilised PMMA [1].  $\beta=1.0$ ,  $\kappa=0.4$  &  $C=3.6$ .

Here the flow depends sensitively on  $\kappa/\beta$  and the balance of friction and adhesion.

Residual adhesion may explain the prolonged shear thinning seen in granular suspensions and their sensitivity.

## III. Strong adhesion: $\sigma_A \gg \sigma^*$



**System:** corn-starch in sunflower oil with lecithin,  $\kappa=1.1$  &  $C=2.3$ .

Attraction alone does not explain the yield stress of some granular suspensions.

The yield stress is instead driven by constraints and adhesion, depending strongly on inter-particle friction.